



Principal of least degeneracy
Higher order matching and nonlinear
equations

Lesson 16

A boundary layer of thickness $O(\sqrt{\varepsilon})$

Let us consider the example $\varepsilon \frac{d^2 f}{dx^2} + x^2 \frac{df}{dx} - f = 0$ with $x \in (0, 1)$, $f(0) = 1$, $f(1) = 2$. We will look for the one term expansion. Since $x^2 > 1$ on $(0, 1)$ we assume that the boundary layer is at $x = 0$. **Is something strange up to now? Thus**

$$x^2 \frac{df_0^{out}}{dx} - f_0^{out} = 0, \quad f_0^{out}(1) = 2.$$

Separating variables leads to

$$\int \frac{df_0^{out}}{f_0^{out}} = \int \frac{dx}{x^2} \Rightarrow f_0^{out} = C e^{-1/x} = 2e^{1-1/x}.$$

We suppose that $s = x/\varepsilon$, then the new equation is

$$\frac{d^2 f}{ds^2} + \varepsilon^2 s^2 \frac{df}{ds} - \varepsilon f = 0, \quad f(0) = 1$$

then

$$\frac{d^2 f_0^{in}}{ds^2} = 0 \Rightarrow f_0^{in} = As + 1.$$

Applying matching condition, we get

$$\lim_{x \rightarrow 0} 2e^{1-1/x} = \lim_{s \rightarrow \infty} As + 1.$$

What is wrong?

Thickness $O(\sqrt{\varepsilon})$

Let us try the stretched variable $s = x/\varepsilon^p$, then

$$\frac{df}{dx} = \frac{1}{\varepsilon^p} \frac{df}{ds}, \quad \frac{d^2 f}{dx^2} = \frac{1}{\varepsilon^{2p}} \frac{d^2 f}{ds^2}.$$

$$\varepsilon^{1-2p} \frac{d^2 f}{ds^2} + \varepsilon^p s^2 \frac{df}{ds} - f = 0.$$

Let us try to keep the maximum number of terms in the equation. Taking $p = 1/2$ we keep the second derivative and f . Then

$$\frac{d^2 f_0^{in}}{ds^2} - f_0^{in} = 0, \quad f_0^{in}(0) = 1, \Rightarrow f_0^{in} = Ae^s + (1 - A)e^{-s}.$$

Applying Prandtl's matching condition, we get

$$\lim_{x \rightarrow 0} 2e^{1-1/x} = \lim_{s \rightarrow \infty} Ae^s + (1-A)e^{-s}$$

that implies $A = 2$. Thus

$$f_0^{comp} = 2e^{1-1/x} + e^{-x/\sqrt{\varepsilon}}.$$

If $p < 1/2$ then $f_0^{in} = 0$ which does not satisfy the boundary condition at $s = 0$.

An interior boundary layer

$$\varepsilon \frac{d^2 f}{dx^2} + x \frac{df}{dx} + x f = 0 \text{ with } x \in (-1, 1), f(-1) = e, f(1) = 2e^{-1}.$$

We expect the boundary layer at $x = 0$, two outer solution $f_0^{+out} = f_0^+$ and $f_0^{-out} = f_0^-$. We have

$$\frac{df_0^+}{dx} + f_0^+ = 0, \quad f_0^+(1) = 2e^{-1} \quad \Rightarrow \quad f_0^+ = 2e^{-x},$$

$$\frac{df_0^-}{dx} + f_0^- = 0, \quad f_0^-(-1) = e \quad \Rightarrow \quad f_0^- = e^{-x}.$$

An interior boundary layer

The stretched variable $s = x/\varepsilon^p$ leads to

$$\varepsilon^{1-2p} \frac{d^2 f}{ds^2} + s \frac{df}{ds} + \varepsilon^p s f = 0.$$

Taking $p = 1/2$, we get

$$\frac{d^2 f_0^{in}}{ds^2} + s \frac{df_0^{in}}{ds} = 0.$$

We make a substitution $w = \frac{df_0^{in}}{ds}$, then

$$\frac{dw}{ds} + sw = 0 \quad \Rightarrow \quad w = Ae^{-s^2/2} \quad \Rightarrow$$

An interior boundary layer

$$f_0^{in} = A \int_0^s e^{-t^2/2} dt + f_0^{in}(0).$$

We write $\int_0^s e^{-t^2/2} dt = \frac{1}{\sqrt{2}} \int_0^{s/\sqrt{2}} e^{-u^2} du = B \operatorname{erf}(s/\sqrt{2})$.

Then

$$f_0^{in} = B \operatorname{erf}(s/\sqrt{2}) + f_0^{in}(0).$$

The Prandtl's condition gives

$$\lim_{x \rightarrow 0^+} 2e^{-x} = \lim_{s \rightarrow +\infty} B \operatorname{erf}(s/\sqrt{2}) + f_0^{in}(0) \Rightarrow B + f_0^{in}(0) = 2,$$

$$\lim_{x \rightarrow 0^-} e^{-x} = \lim_{s \rightarrow +\infty} B \operatorname{erf}(s/\sqrt{2}) + f_0^{in}(0) \Rightarrow -B + f_0^{in}(0) = 1.$$

An interior boundary layer

$$B = 1/2, \quad f_0^{in}(0) = 3/2.$$

Summarizing

$$f_0^+ = 2e^{-x}, \quad x > O(\sqrt{\varepsilon})$$

$$f_0^{in} = 1/2 \operatorname{erf}(x/\sqrt{2\varepsilon}) + 3/2, \quad x = O(\sqrt{\varepsilon})$$

$$f_0^- = e^{-x}, \quad x < -O(\sqrt{\varepsilon})$$

The composition solution can not be produced by the standard way. We observe

An interior boundary layer

$$f_0^{in}(x > O(\sqrt{\varepsilon})) = 1/2 + 3/2 + \text{sm. t.}, \text{ since } \text{erf} \sim 1,$$

$$f_0^{in}(x < -O(\sqrt{\varepsilon})) = -1/2 + 3/2 + \text{sm. t.}, \text{ since } \text{erf} \sim -1,$$

We conclude that

$$f_0^{comp} = \left(1/2 \text{erf}(x/\sqrt{2\varepsilon}) + 3/2 \right) e^{-x}.$$

An interior boundary layer

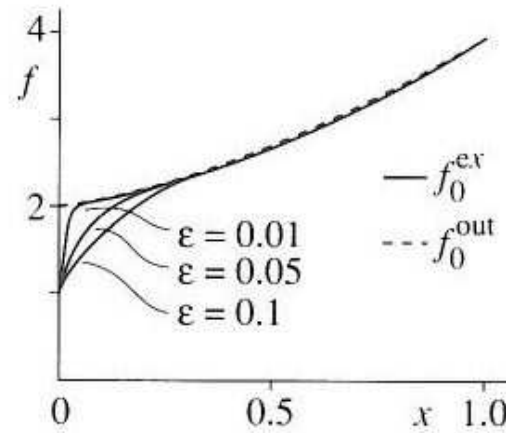


Fig. 5.1 The exact solution for various values of ϵ

Higher order matching

We shall work with the old example

$$\varepsilon \frac{d^2 f}{dx^2} + \frac{df}{dx} = 2x + 1, \quad x \in (0, 1), \quad f(0) = 1, \quad f(1) = 4,$$

$$f^{ex} = x^2 + x + 2 - e^{-x/\varepsilon} + \varepsilon \left(2(1 - x) - 2e^{-x/\varepsilon} \right)$$

Using the stretched variable $s = x/\varepsilon$ we get

Higher order matching

$$\frac{d^2 f}{ds^2} + \frac{df}{ds} = 2\varepsilon^2 s + \varepsilon, \quad f(s, \varepsilon) \sim f_0(s) + \varepsilon f_1(s) + \dots,$$

$$\varepsilon^0 \frac{d^2 f_0^{in}}{ds^2} + \frac{df_0^{in}}{ds} = 0, \quad f_0^{in}(0) = 1 \Rightarrow f_0^{in} = A + (1 - A)e^{-s}$$

$$\varepsilon^1 \frac{d^2 f_1^{in}}{ds^2} + \frac{df_1^{in}}{ds} = 1, \quad f_1^{in}(0) = 0 \Rightarrow f_1^{in} = B - Be^{-s} + s$$

$$\varepsilon^2 \frac{d^2 f_2^{in}}{ds^2} + \frac{df_2^{in}}{ds} = 2s, \quad f_2^{in}(0) = 0 \Rightarrow f_2^{in} = C - Ce^{-s} + s^2 - 2s$$

$$\varepsilon^n \frac{d^2 f_n^{in}}{ds^2} + \frac{df_n^{in}}{ds} = 0, \quad f_n^{in}(0) = 0 \Rightarrow f_n^{in} = D_n - D_n e^{-s},$$

$$n = 3, 4, \dots$$

Higher order matching

$$f^{in} = A + (1 - A)e^{-s} + \varepsilon(B - Be^{-s} + s) \\ + \varepsilon^2(C - Ce^{-s} + s^2 - 2s) + \sum_{n=3}^{\infty} \varepsilon^n D_n(1 - e^{-s}).$$

The Prandtl's condition can not be applied, because of existence of unbounded terms εs and $\varepsilon^2(s^2 - 2s)$. They represent x and $x^2 - 2\varepsilon x$. Let us write the exact solution in terms of s :

Higher order matching

$$\begin{aligned} f^{ex} &= x^2 + x + 2 - e^{-x/\varepsilon} + \varepsilon \left(2(1 - x) - 2e^{-x/\varepsilon} \right) \\ &= \varepsilon^2 s^2 + \varepsilon s + 2 - e^{-s} + \varepsilon \left(2(1 - \varepsilon s) - 2e^{-s} \right) \\ &= 2 - e^{-s} + \varepsilon(2 - 2e^{-s} + s) + \varepsilon^2(s^2 - 2s). \end{aligned}$$

Comparing with f^{in} we get $A = 2, B = 2, C = 0, D_n = 0, n = 3, 4, \dots$. How we can define A, B, C, D_n without of exact solution?

Higher order matching

Let us look on the Prandtl's condition from the other point of view. Let $x = O(\varepsilon^{1/2})$, then

$$f^{in} = A + (1-A)e^{-s} = A + o(1), \quad f^{out} = x^2 + x + 2 - e^{-x/\varepsilon} = 2 + o(1)$$

We recuperated $A = 2$. Let us do the same, but in slightly general form. We represent $s = \frac{x}{\varepsilon}$ in two variables $x = t\varepsilon^\alpha$ then $s = \frac{t}{\varepsilon^{1-\alpha}}$. Let us represent two term expansion in terms of new variable $t = x/\varepsilon^\alpha$, $\alpha \in (0, 1)$ for t fixed. We get

$$f^{out} = x^2 + x + 2 + \varepsilon 2(1 - x) = 2 + \varepsilon^\alpha t + \varepsilon^{2\alpha} t^2 + 2\varepsilon - 2\varepsilon^{1+\alpha} t$$

Higher order matching

Now we use $s = \frac{t}{\varepsilon^{1-\alpha}}$ and find two term expansion of f^{in} in terms of $\frac{t}{\varepsilon^{1-\alpha}}$.

$$f^{in} = A + \varepsilon(B + s) = A + \varepsilon^\alpha t + \varepsilon B.$$

Matching with

$$f^{out} = 2 + \varepsilon^\alpha t + \varepsilon^{2\alpha} t^2 + 2\varepsilon - 2\varepsilon^{1+\alpha} t$$

we have to ask $\varepsilon^{2\alpha} = o(\varepsilon)$, or $\alpha > 1/2$. We get $A = 2$ and $B = 2$.

Higher order matching

The composition solutions is

$$f_{2term}^{comp} = f_{2term}^{out} + f_{2term}^{in} - f_{2term}^{match}$$

where f_{2term}^{comp} is given by terms of order up to $O(\varepsilon)$. In our case it is $2 + x + 2\varepsilon$. We have

$$f^{comp} = x^2 + x + 2 + 2\varepsilon(1 - x)$$

$$+ 2 - e^{-x/\varepsilon} + \varepsilon(2 - 2e^{-x/\varepsilon} + x/\varepsilon)$$

$$-2 - x - 2\varepsilon = x^2 + x + 2 - e^{-x/\varepsilon} + \varepsilon(2 - 2x - 2e^{-x/\varepsilon}).$$

Nonlinear examples

$$\varepsilon \frac{d^2 f}{dx^2} + \frac{df}{dx} + f^2 = 0, \quad x \in (0, 1), f(0) = 2, f(1) = 1/2.$$

$$\frac{df_0^{\text{out}}}{dx} + (f_0^{\text{out}})^2 = 0, \quad f_0^{\text{out}}(1) = 1/2 \Rightarrow f_0^{\text{out}} = \frac{1}{x+1}.$$

$$s = \frac{x}{\varepsilon} \Rightarrow \frac{d^2 f^{\text{in}}}{ds^2} + \frac{df^{\text{in}}}{ds} + \varepsilon (f^{\text{in}})^2 = 0$$

$$\frac{d^2 f_0^{\text{in}}}{ds^2} + \frac{df_0^{\text{in}}}{ds} = 0 \Rightarrow f_0^{\text{in}} = A + (2 - A)e^{-s}$$

The Prandtl's condition implies

$$\lim_{x \rightarrow 0} \frac{1}{x+1} = \lim_{s \rightarrow \infty} A + (2-A)e^{-s} \Rightarrow A = 1.$$

$$f^{comp} = \frac{1}{x+1} + e^{-x/\varepsilon}.$$

Nonlinear examples

$$\varepsilon \frac{d^2 f}{dx^2} + 2f \frac{df}{dx} - 4f = 0, \quad x \in (0, 1), f(0) = 0, f(1) = 4.$$

$$2f_0^{out} \frac{df_0^{out}}{dx} - 4f_0^{out} = 0, f(1) = 4 \Rightarrow f_0^{out} = 2x + 2$$

$$s = \frac{x}{\varepsilon} \Rightarrow \frac{d^2 f_0^{in}}{ds^2} + 2f_0^{in} \frac{df_0^{in}}{ds} = 0 \Rightarrow \frac{d}{ds} \left(\frac{df_0^{in}}{ds} + (f_0^{in})^2 \right) = 0$$

$$\frac{df_0^{in}}{ds} + (f_0^{in})^2 = c = a^2$$

since $f_0^{in}(0) = 0$ and $\lim_{x \rightarrow 0} f_0^{out} = 2$ and f_0^{in} increase monotonically.

Nonlinear examples

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- $A = 1, f_0^{in} = a \frac{e^{2as} - 1}{e^{2as} + 1} = a \tanh as$
- $\lim_{x \rightarrow 0} f_0^{out} = \lim_{s \rightarrow \infty} f_0^{in} \Rightarrow a = 2$

Nonlinear examples

$$f_0^{in} = a \tanh 2s$$

$$f^{comp} = 2x + 2 \tanh(2x/\varepsilon).$$

The end